

Strong Enhancement of the Critical Current at the Antiferromagnetic Transition in $\text{ErNi}_2\text{B}_2\text{C}$ Single Crystals

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We report on transport and magnetization measurements of the critical current density J_c in $\text{ErNi}_2\text{B}_2\text{C}$ single crystals that show strongly enhanced vortex pinning at the Néel temperature T_N and low applied fields. The height of the observed J_c peak decreases with increasing magnetic field in clear contrast with that of the peak effect found at the upper critical field. We also performed the first angular transport measurements of J_c ever conducted on this compound. They reveal the correlated nature of this pinning enhancement, which we attribute to the formation of antiphase boundaries at T_N .

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Single quanta of magnetic flux enter a type II superconductor in the form of vortices when it is exposed to a magnetic field H . This is a big setback for applications, because when an electrical current J is applied, vortices move and energy is dissipated. This movement can be arrested by non-superconducting defects that pin vortices by lowering the system energy. The interaction between vortices and pinning centers has been studied extensively for decades, focused mainly on the so-called vortex core pinning (caused by the local suppression of the superconducting order parameter) [1, 2]. Less explored is the interplay between vortices and magnetic media, which could offer pinning forces superior to those from core pinning [3–6]. A common difficulty of studying magnetic pinning is the inability of separating the magnetic and core contributions. It is, therefore, of great advantage to study systems where the magnetic transition takes place inside the superconducting phase, allowing one to compare the behavior with and without the magnetically ordered phase.

There are a considerable number of materials in which to study the coexistence between superconductivity and ordered magnetic phases, such as the Chevrel-phases [7], CeCoIn_5 [8], and recently iron pnictides [9–11]. However, the rare earth-nickel-borocarbide family ($\text{RENi}_2\text{B}_2\text{C}$, where RE is a rare earth element) [12–14] has several advantages, namely a relatively high superconducting transition temperature T_c and the tunability of the ratio between its magnetic and superconducting ordering temperatures, which can be changed by using different rare earth elements [15, 16]. The readily available high-purity single crystals allow one to study the effects of magnetic pinning without any significant defect (non-magnetic) contribution [15].

Among the borocarbide family the compound with $\text{RE} = \text{Er}$ has a T_c of 10.5 K and a Néel temperature T_N of 6.0 K [17]. The occurrence of antiferromagnetism directly influences the superconducting properties, as seen in the upper critical field vs temperature curve, where $H_{c2}(T)$ is

slightly suppressed just below $T = T_N$ for $\mathbf{H} \parallel c$ and has an inflection point for $\mathbf{H} \parallel ab$ [17]. This is a consequence of the local moments ordering in the antiferromagnetic phase, as was shown experimentally [17, 18] and theoretically [19, 20].

In addition to the occurrence of antiferromagnetism, at $T = T_N$ a tetragonal to orthorhombic crystal structure transition takes place [21]. The resulting twin boundaries were shown by Bitter decoration, magneto-optical, and scanning Hall-probe experiments to act as pinning centers at lower temperatures ($T < T_N$) [22, 23]. The pinning mechanism was proposed to be caused by a ferromagnetic spin component parallel to the crystallographic c axis localized at the twin boundaries [22].

Although borocarbides (and particularly the compound with $\text{RE} = \text{Er}$) have been extensively studied, few critical current measurements by electrical transport have been performed, with the exception of the seminal work by Gammel *et al.* [24], exploring the weak ferromagnetism below $T = 2.3$ K [25–28]. This was shown to lead to an enhanced critical current density J_c , which was ascribed to an increase in pinning force due to pair breaking by the ferromagnetism [24]. Canfield *et al.* [29] pointed out that the data in Ref. [24] extrapolated to $J_c = 0$ at T_N , indicating a clear linkage between pinning and the antiferromagnetic order. However, this extrapolation could not be confirmed because transport experiments were only done for $T < 4.2$ K to take advantage of the cooling power of liquid helium [24].

The lack of transport experiments is to a great extent due to heating problems caused by the applied current being very high because of the large cross section of single crystals. To the best of our knowledge, the present work is the first report both of transport measurements on $\text{ErNi}_2\text{B}_2\text{C}$ single crystals at temperatures around T_N and of angular J_c measurements performed on this compound.

In this Letter, we present the results of transport and magnetization measurements on $\text{ErNi}_2\text{B}_2\text{C}$ single

crystals, which revealed a local maximum in $J_c(T)$ at $T = T_N$. We study this large increase in J_c as a function of field strength and orientation, and we find that it occurs only for fields oriented along the crystallographic c axis. Unlike the peak effect near H_{c2} , the height of the newly discovered maximum decreases with increasing magnetic field. We rule out pinning from twin boundaries as we are able to determine its angular fingerprint at lower temperatures, which differs from the increase in J_c observed at T_N . We attribute the findings to vortex pinning due to antiphase boundaries between antiferromagnetic domains.

Large, homogeneous single crystals of $\text{ErNi}_2\text{B}_2\text{C}$ were grown using the Ni_2B flux growth technique [30]. The samples used for transport measurements were polished mechanically down to a thickness of 40 – 60 μm parallel to the c axis before being cut to pieces ~ 250 μm wide and ~ 1.5 mm long. Sputtered Au contacts allowed us to apply currents along the ab -planes while measuring the voltage across the sample. We observed a slight reduction of T_c (~ 0.5 K) due to sample preparation, which is consistent with the previously reported increase in T_c after annealing [31]. The small decrease in T_c does not affect the overall results of the present study; it left both T_N and the shape of $H_{c2}(T)$ unchanged.

Transport measurements were performed in three systems characterized by different cooling methods: (i) a Quantum Design, Inc. PPMS with a semi-closed He system, (ii) a variable temperature insert with flowing He, and (iii) with the sample immersed in liquid He. In all systems we used high-precision rotators and the maximum Lorentz force configuration ($\mathbf{J} \perp \mathbf{H}$). The critical current density was obtained from IV -curves using a 1 μV criterion. Measurements of $R(T)$ and $R(H)$ gave the upper critical field vs temperature, using a value of R corresponding to 90% of the normal conducting resistance. The comparison of the J_c results obtained in the three systems allowed us to assess the magnitude of heating effects.

Magnetization measurements were carried out in a Quantum Design, Inc. MPMS Superconducting Quantum Interference Device (SQUID). J_c was obtained using the Bean critical state model with $J_c = \frac{|m_i|}{abc} \cdot \frac{4}{b(1-b/3a)}$, where $a \geq b$, m_i is the irreversible magnetic moment, and a , b , and c are the sample length, width, and thickness, respectively [32].

In Fig. 1(a) we present a J_c transport measurement vs temperature at $\mu_0 H = 0.5$ T for $\mathbf{H} \parallel c$. Starting at low temperature, $J_c(T)$ decreases rapidly as T increases, and a rough extrapolation of $J_c(T)$ leads to 0 at $T = T_N$, the same behavior found in Ref. [24]. Our new measurements presented here, however, show that J_c instead changes abruptly at T_N . The slope of $J_c(T)$ is much lower at $T > T_N$ than below, and a peak in $J_c(T)$ is observed at T_N , corresponding to a two-fold increase compared to the

value just below and above it. These are the key results of this Letter. Finally, at a temperature just above 7 K we observe the well-documented peak-effect at H_{c2} .

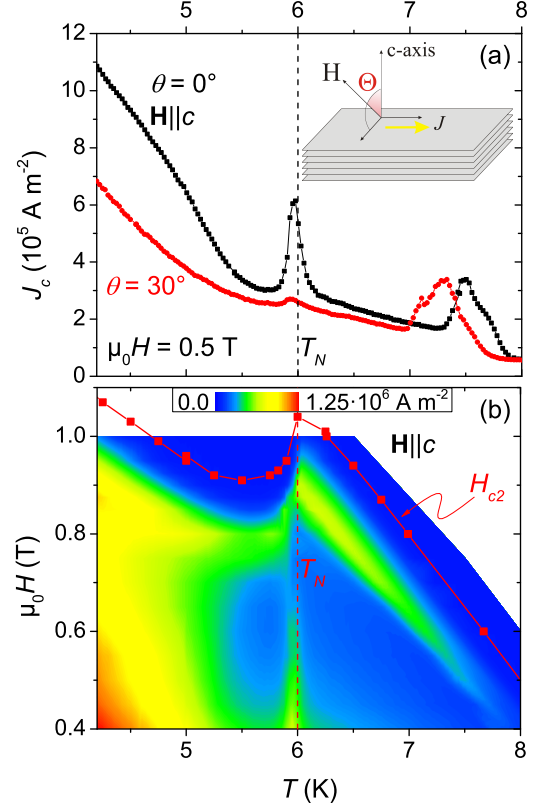


FIG. 1. (Color online) (a) Transport measurements of J_c for $\mathbf{H} \parallel c$ and $\theta = 30^\circ$. A new peak in J_c is observed at $T = T_N$ for $\mathbf{H} \parallel c$. The sketch gives the measurement geometry. (b) Contour plot of J_c as a function of temperature and magnetic field $\mathbf{H} \parallel c$ showing the field range where the peak at $T = T_N$ is present.

The maximum at $T = T_N$ occurs over a range of fields, as seen in the contour plot in Fig. 1(b), obtained from multiple $J_c(T)$ measurements at different $\mathbf{H} \parallel c$, where different colors represent different values of J_c . Note that the height of the J_c peak at T_N shows the opposite field dependence than that of the one observed near H_{c2} , indicating a different origin of the two maxima. Whereas the peak near H_{c2} becomes more pronounced as H is increased, the peak at T_N decreases with increasing H until it has essentially vanished at $\mu_0 H = 0.7$ T. This leads us to speculate that the peak at T_N has a magnetic origin, as magnetic pinning tends to be attenuated as the magnetic field modulation of the vortices becomes smaller. Our observations are consistent with the dip found in the dynamic magnetic susceptibility of $\text{ErNi}_2\text{B}_2\text{C}$ single crystals at $T = T_N$, measured by Prozorov *et al.* [33] using a tunnel-diode resonator technique. They explained their results by a pinning enhancement due to the occurrence of antiferromagnetic order, accompanied by large

magnetic fluctuations. The decrease in the peak height with increasing H also rules out the drastic change in $H_{c2}(T)$ at T_N as the origin for the maximum in $J_c(T)$, since the peak disappears far below $H_{c2}(T_N)$. The absence of a $J_c(T)$ peak at $\theta = 30^\circ$, shown in Fig. 1(a), is further proof that this peak is not associated with the sudden change in $H_{c2}(T)$ at T_N , since for $\theta = 30^\circ$ this feature in $H_{c2}(T)$ is still visible (not shown).

It is clear from the $J_c(T)$ measurements at $\theta = 0^\circ$ and 30° that this effect is strongly angle-dependent. Thus, further insight can be gained from the dependence of J_c on field orientation, since it allows one to distinguish between the superconducting and the magnetic anisotropy.

The obtained angular J_c curves are reproduced in Fig. 2(a). They reveal rich and complex features. Despite the small change of the absolute value of H_{c2} with field orientation, we found that J_c has a pronounced maximum for $\mathbf{H} \parallel ab$ at all temperatures investigated. This indicates strong pinning for $\mathbf{H} \parallel ab$, in contrast to what would be expected from the anisotropy of H_{c2} , where $H_{c2}^{\parallel c} > H_{c2}^{\parallel ab}$. It also differs from what has been observed in non-magnetic $\text{YNi}_2\text{B}_2\text{C}$ thin films, where both J_c and H_{c2} show the same anisotropy [34, 35]. The maximum we find in $J_c(\theta)$ around $\mathbf{H} \parallel ab$ can be explained by magnetic pinning, and it is consistent with the angular dependence of the magnetic susceptibility of the Er ions, namely $\chi \sim 0$ for $\mathbf{H} \parallel c$, while for $\mathbf{H} \parallel ab$ the full Curie-Weiss form is observed [16]. We also find small maxima at $\theta = \pm 22^\circ$ both above and below T_N , which are not related to the sample shape and will require further investigation.

At $T = T_N$ a wide peak centered at $\mathbf{H} \parallel c$ is visible, whereas at temperatures only 0.25 K above and below a plateau is observed. The peak in the angular dependence of J_c indicates that the nature of this pinning is correlated. Naturally one is tempted to ascribe this maximum to the already reported pinning by a ferromagnetic spin component localized at magnetic twin boundaries [22]; however we have measured a much narrower $J_c(\theta)$ peak at lower temperatures [$T = 4.2$ K, see the inset of Fig. 2(a)], which is clearly due to the presence of twin boundaries, as has been observed in Bitter decoration and scanning Hall-probe experiments [22, 23]. Its different shape indicates that the J_c peak at T_N is *not* caused by twin boundaries. The peak due to twin boundaries appears for $T < T_N$ and exhibits a monotonic growth as the temperature decreases. The height of this peak can be roughly extracted when comparing $J_c(T)$ at $\theta = 0^\circ$ and 30° shown in Fig. 1(a). The disappearance of this peak just below T_N is also confirmed by the angular dependence at $T = 5.75$ K [see Fig. 2(a)]. The temperature evolution of $J_c(\theta)$ indicates that the peak at T_N has a different origin than that at low temperatures, although both are correlated along the c axis.

The dependence of the T_N peak on temperature and field orientation becomes even clearer in the contour plot

in Fig. 2(b). It can be seen that J_c is enhanced only over a narrow temperature and angular range near $T = 6$ K and $\theta = 0^\circ$.

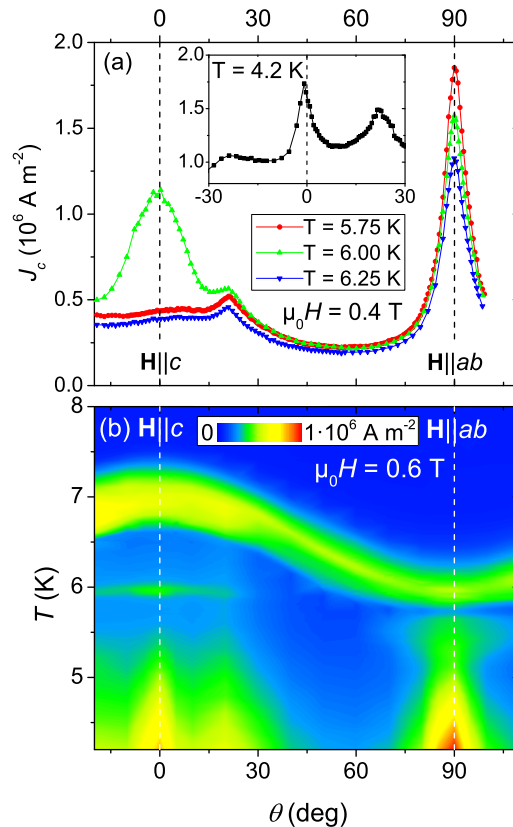


FIG. 2. (Color online) (a) Angular transport J_c measurements at $\mu_0 H = 0.4$ T reveal the narrow temperature range at T_N over which the peak at $\mathbf{H} \parallel c$ occurs. (b) The contour plot of J_c as a function of magnetic field orientation θ and T shows the location of the peak at T_N and $\theta = 0^\circ$ in the T - θ space.

The correlated nature of the pinning and its location around T_N lead us to conclude that it is due to dynamic magnetic domains that are created at T_N and annihilated as the temperature decreases further. Domain walls are known to produce planar pinning potentials in ferromagnetic/superconducting hybrids [36, 37]. This poses the interesting question why the pinning enhancement occurs for $\mathbf{H} \parallel c$, while the easy axis of the magnetic moment for $\text{ErNi}_2\text{B}_2\text{C}$ is the b axis [38, 39]. This could be due to the magnetic moment being flipped to the c axis; it has to be noted, however, that this is highly unlikely because of the high crystalline electric field anisotropy of the local moment sublattice. Another explanation could be vortices bending into the ab planes at the magnetic domain boundaries in order to gain the Zeeman energy. As the vortices move away from the c axis the energy gain gets gradually smaller, until at the ab planes there is no gain since vortices have zero net gain. A phenomenological model indicates that $J_c \propto \cos^3(\theta)$ [40]. This

gives a broader maximum than the one observed experimentally, nevertheless it can explain qualitatively the enhancement observed. The presented results and the proposed model motivate measurements in other compounds such as $\text{TmNi}_2\text{B}_2\text{C}$, which has its easy axis parallel to the c axis [41]. It is worth noting that associated with the tetragonal to orthorhombic transition in $\text{ErNi}_2\text{B}_2\text{C}$ [21], strain occurs that could also be responsible for enhanced pinning [42]. Whether this is indeed the case could be answered by investigating a borocarbide which does not undergo a structural transition at $T = T_N$, namely again $\text{TmNi}_2\text{B}_2\text{C}$.

In order to further explore the effects at lower fields, which are not accessible to transport measurements due to heating, different samples were taken from the same batch and investigated in the SQUID. This also serves as verification of our transport measurement results. After creating a Bean profile [32] by applying a field $\mathbf{H} \parallel c$ larger than $2H^*$ (the field above which screening currents flow in the entire sample volume) at $T = 4.5$ K, the magnetic moment m was measured vs temperature at $\mu_0 H = 0.02$ T. In order to capture the increase in J_c around T_N the Bean profile was generated at different temperatures near T_N before setting $\mu_0 H = 0.02$ T and measuring $m(T)$.

We then compiled an $m(T)$ curve from the different branches, selecting the branch with the highest (absolute) value of m for each temperature. The obtained curve represents an envelope of all $m(T)$ branches. With the Bean model the irreversible magnetic moment then gives the critical current density [32].

As can be seen in Fig. 3, the obtained $J_c(T)$ shows a peak exactly at $T = T_N$, similar to the local maximum observed in the transport measurements [see Fig. 1(a)]. Although the shape of $J_c(T)$ at this field resembles that of $H_{c2}(T)$, it is not directly governed by $H_{c2}(T)$, as the height of the J_c peak obtained from magnetization curves decreases with increasing magnetic field, vanishing far below $H_{c2}(T_N)$, as in transport measurements. The peak in J_c at T_N and its field dependence are seen in magnetic measurements for both a sample cut and polished to similar dimensions as the one used in transport experiments and a pristine single crystal with no treatment at all. This reinforces that the phenomenon is not sample dependent and that it can be experimentally verified by different techniques.

In summary, we have observed in transport and magnetization measurements an enhancement of the critical current density in $\text{ErNi}_2\text{B}_2\text{C}$ single crystals at the Néel temperature. By performing the first angular transport measurements we determined that this increase occurs for magnetic fields applied parallel to the c axis, consistent with vortex pinning due to antiphase boundaries between antiferromagnetic domains. This study opens a new avenue to investigate the interaction between superconductivity and different magnetic phases, such as those

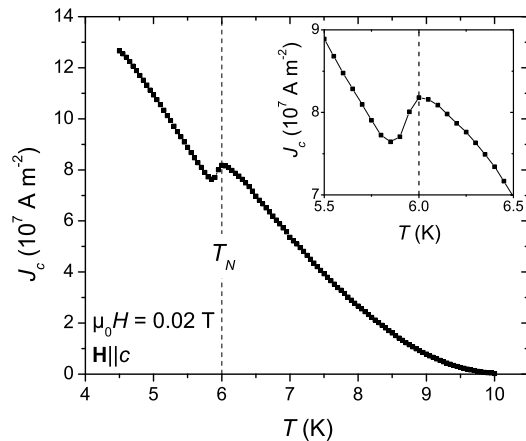


FIG. 3. The critical current density vs temperature at $\mu_0 H = 0.02$ T (with $\mathbf{H} \parallel c$), obtained from magnetization measurements, exhibits a local maximum at $T = T_N$.

known in Er, Ho and Tm borocarbides. In addition, new developments can be expected for other superconductors with coexistence of magnetic phases, such as underdoped copper- and iron-based superconductors [43, 44].

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